

SOME USEFUL FORMULAE IN $SU(N)$ ALGEBRA*

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ABSTRACT

Some useful relations in $SU(N)$ algebra are given in this paper. These relations may be useful in deriving the many-body interaction in QCD.

Keywords $SU(N)$ algebra, Many body interaction, Quantum chromodynamics, Hermitian operators, Many-dimensional calculations

1 INTRODUCTION

In particle physics, $SU(N)$ algebra is an useful tool. For example, $SU(2)$ is the basis of discussing the angular momentum, isospin and the weak-electromagnetic unification; $SU(3)$ is the basis of discussing QCD and the flavour group of light quark; $SU(4)$ is the basis of Dirac algebra; $\dots$, and so on. In this paper, we follow the way introduced in Ref.[1] to give some useful relations in $SU(N)$ algebra. Consider a set of $N^2 - 1$ traceless Hermitian operator $\lambda_k (k = 1, 2, \dots, N^2 - 1)$ defined a N -dimensional space.

$$\lambda_k^\dagger = \lambda_k \quad (1)$$

$$Tr(\lambda_k) = 0 \quad (2)$$

$$\lambda_k \lambda_l = \alpha \delta_{kl} I + \beta_{klm} \lambda_m \quad (3)$$

$$\beta_{klm} = d_{klm} + i f_{klm} \quad (4)$$

$$[\lambda_k, \lambda_l] = \lambda_k \lambda_l - \lambda_l \lambda_k = 2i f_{klm} \lambda_m \quad (5)$$

$$\{\lambda_k, \lambda_l\} = \lambda_k \lambda_l + \lambda_l \lambda_k = 2\alpha \delta_{kl} + 2d_{klm} \lambda_m \quad (6)$$

Consider a set of column vectors; $\zeta_1, \zeta_2, \dots$; and a set of row vectors; $\eta_1, \eta_2, \dots$, $N \times N$ matrix $\zeta_1 \eta_1$, $\zeta_1 \eta_1 \zeta_2 \eta_2$, and $\zeta_1 \eta_1 \zeta_2 \eta_2 \zeta_3 \eta_3$ can be written as:

$$\zeta_1 \eta_1 = \frac{1}{N} (\eta_1 \zeta_1) + \frac{1}{N\alpha} (\eta_1 \lambda_k \zeta_1) \lambda_k \quad (7)$$

$$\begin{aligned} \zeta_1 \eta_1 \zeta_2 \eta_2 &= \frac{1}{N^2} (\eta_1 \zeta_1) (\eta_2 \zeta_2) + \frac{1}{N^2 \alpha} (\eta_1 \zeta_1) (\eta_2 \lambda_l \zeta_2) \lambda_l \\ &+ \frac{1}{N^2 \alpha} (\eta_1 \lambda_k \zeta_1) (\eta_2 \zeta_2) \lambda_k + \frac{1}{N^2 \alpha^2} (\eta_1 \lambda_k \zeta_1) (\eta_2 \lambda_l \zeta_1) \lambda_k \lambda_l \end{aligned} \quad (8)$$

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$$\begin{aligned}
& \zeta_1 \eta_1 \zeta_2 \eta_2 \zeta_3 \eta_3 \\
= & \frac{1}{N^3} (\eta_1 \zeta_1) (\eta_2 \zeta_2) (\eta_3 \zeta_3) + \frac{1}{N^3 \alpha} (\eta_1 \zeta_1) (\eta_2 \lambda_l \zeta_2) (\eta_3 \zeta_3) \lambda_l \\
& + \frac{1}{N^3 \alpha} (\eta_1 \lambda_k \zeta_1) (\eta_2 \zeta_2) (\eta_3 \zeta_3) \lambda_k + \frac{1}{N^3 \alpha^2} (\eta_1 \lambda_k \zeta_1) (\eta_2 \lambda_l \zeta_2) (\eta_3 \zeta_3) \lambda_k \lambda_l \\
& + \frac{1}{N^3 \alpha} (\eta_1 \zeta_1) (\eta_2 \zeta_2) (\eta_3 \lambda_m \zeta_3) \lambda_m + \frac{1}{N^3 \alpha^2} (\eta_1 \zeta_1) (\eta_2 \lambda_l \zeta_2) (\eta_3 \lambda_m \zeta_3) \lambda_l \lambda_m \\
& + \frac{1}{N^3 \alpha^2} (\eta_1 \lambda_k \zeta_1) (\eta_2 \zeta_2) (\eta_3 \lambda_m \zeta_3) \lambda_k \lambda_m + \frac{1}{N^3 \alpha^3} (\eta_1 \lambda_k \zeta_1) (\eta_2 \lambda_l \zeta_2) (\eta_3 \lambda_m \zeta_3) \lambda_k \lambda_l \lambda_m \quad (9)
\end{aligned}$$

All these relations can be found in Ref.[1] except Eq.(9), which can be derived by the same technique used in Ref.[1] to obtain the relation . A summation conversion of repeated indices is assumed here.

2 DERIVATION

In order to give the useful formulae, we also need the following relations(see Ref.[1]):

$$(\lambda_k)_{a_1 b_1} (\lambda_k)_{a_2 b_2} = \alpha (N \delta_{a_1 b_2} \delta_{a_2 b_1} - \delta_{a_1 b_1} \delta_{a_2 b_2}) \quad (10)$$

$$\begin{aligned}
& \beta_{klm} (\lambda_k)_{a_1 b_1} (\lambda_l)_{a_2 b_2} (\lambda_m)_{a_3 b_3} \\
= & \alpha^2 [N^2 \delta_{a_1 b_2} \delta_{a_2 b_3} \delta_{a_3 b_1} - N (\delta_{a_1 b_1} \delta_{a_2 b_3} \delta_{a_3 b_2} + \delta_{a_1 b_3} \delta_{a_2 b_2} \delta_{a_3 b_1} \\
& + \delta_{a_1 b_2} \delta_{a_2 b_1} \delta_{a_3 b_3}) + 2 \delta_{a_1 b_1} \delta_{a_2 b_2} \delta_{a_3 b_3}] \quad (11)
\end{aligned}$$

Consider further the following $N \times N$ matrix;

$$\begin{aligned}
& (\eta_4 \zeta_1) (\eta_1 \zeta_2) (\eta_2 \zeta_3) (\eta_3 \zeta_4) \\
= & \frac{1}{N^3} (\eta_1 \zeta_1) (\eta_2 \zeta_2) (\eta_3 \zeta_3) (\eta_4 \zeta_4) + \frac{1}{N^3 \alpha} (\eta_1 \zeta_1) (\eta_2 \lambda_l \zeta_2) (\eta_3 \zeta_3) (\eta_4 \lambda_l \zeta_4) \\
& + \frac{1}{N^3 \alpha} (\eta_1 \lambda_k \zeta_1) (\eta_2 \zeta_2) (\eta_3 \zeta_3) (\eta_4 \lambda_k \zeta_4) + \frac{1}{N^3 \alpha^2} (\eta_1 \lambda_k \zeta_1) (\eta_2 \lambda_l \zeta_2) (\eta_3 \zeta_3) (\eta_4 \lambda_k \lambda_l \zeta_4) \\
& + \frac{1}{N^3 \alpha} (\eta_1 \zeta_1) (\eta_2 \zeta_2) (\eta_3 \lambda_m \zeta_3) (\eta_4 \lambda_m \zeta_4) + \frac{1}{N^3 \alpha^2} (\eta_1 \zeta_1) (\eta_2 \lambda_l \zeta_2) (\eta_3 \lambda_m \zeta_3) (\eta_4 \lambda_l \lambda_m \zeta_4) \\
& + \frac{1}{N^3 \alpha^2} (\eta_1 \lambda_k \zeta_1) (\eta_2 \zeta_2) (\eta_3 \lambda_m \zeta_3) (\eta_4 \lambda_k \lambda_m \zeta_4) \\
& + \frac{1}{N^3 \alpha^3} (\eta_1 \lambda_k \zeta_1) (\eta_2 \lambda_l \zeta_2) (\eta_3 \lambda_m \zeta_3) (\eta_4 \lambda_k \lambda_l \lambda_m \zeta_4) \quad (12)
\end{aligned}$$

where

$$\begin{aligned}
(\eta_4 \lambda_k \lambda_l \lambda_m \zeta_4) &= \eta_4 (\alpha \delta_{kl} I + \beta_{kl n} \lambda_n) \lambda_m \eta_4 \\
&= \alpha \delta_{kl} \eta_4 \lambda_m \zeta_4 + \beta_{kl n} \eta_4 \lambda_n \lambda_m \zeta_4 \\
&= \alpha \delta_{kl} \eta_4 \lambda_m \zeta_4 + \beta_{kl n} \eta_4 (\alpha \delta_{nm} + \beta_{nmo} \lambda_o) \zeta_4 \\
&= \alpha \delta_{kl} (\eta_4 \lambda_m \zeta_4) + \alpha \delta_{nm} \beta_{kl n} (\eta_4 \zeta_4) + \beta_{kl n} \beta_{nmo} (\eta_4 \lambda_o \zeta_4) \quad (13)
\end{aligned}$$

Substituting Eq.(11) into Eq.(12) and Eq.(10) into Eq.(11), after a tedious calculation, one has seen:

$$\begin{aligned}
 & \beta_{kln}\beta_{nmo}(\lambda_k)_{a_1b_1}(\lambda_l)_{a_2b_2}(\lambda_m)_{a_3b_3}(\lambda_o)_{a_4b_4} \\
 = & \alpha^3(N^3\delta_{a_1b_2}\delta_{a_2b_3}\delta_{a_3b_4}\delta_{a_4b_1} - N^2\delta_{a_1b_2}\delta_{a_2b_4}\delta_{a_3b_3}\delta_{a_4b_1} - N^2\delta_{a_1b_3}\delta_{a_2b_2}\delta_{a_3b_4}\delta_{a_4b_1} \\
 & - N^2\delta_{a_1b_1}\delta_{a_2b_3}\delta_{a_3b_4}\delta_{a_4b_2} - N^2\delta_{a_1b_2}\delta_{a_2b_1}\delta_{a_3b_4}\delta_{a_4b_3} - N^2\delta_{a_1b_2}\delta_{a_2b_3}\delta_{a_3b_1}\delta_{a_4b_4} \\
 & + 2N\delta_{a_1b_1}\delta_{a_2b_2}\delta_{a_3b_4}\delta_{a_4b_3} + 2N\delta_{a_1b_2}\delta_{a_2b_1}\delta_{a_3b_3}\delta_{a_4b_4} + N\delta_{a_1b_4}\delta_{a_2b_2}\delta_{a_3b_3}\delta_{a_4b_1} \\
 & + N\delta_{a_1b_1}\delta_{a_2b_4}\delta_{a_3b_3}\delta_{a_4b_2} + N\delta_{a_1b_1}\delta_{a_2b_3}\delta_{a_3b_2}\delta_{a_4b_4} + N\delta_{a_1b_3}\delta_{a_2b_2}\delta_{a_3b_1}\delta_{a_4b_4} \\
 & - 4\delta_{a_1b_1}\delta_{a_2b_2}\delta_{a_3b_3}\delta_{a_4b_4}) \tag{14}
 \end{aligned}$$

Setting $b_2 = a_4$ and sum over a_4 , one has

$$\begin{aligned}
 & \beta_{kln}\beta_{nmo}(\lambda_k)_{a_1b_1}(\lambda_l)_{a_2b_2}(\lambda_o\lambda_m)_{a_2b_4} \\
 = & \alpha^3[-2N^2\delta_{a_1b_3}\delta_{a_3b_4}\delta_{a_2b_1} - N^2\delta_{a_1b_4}\delta_{a_2b_3}\delta_{a_3b_1} + 3N^2\delta_{a_1b_1}\delta_{a_2b_3}\delta_{a_3b_4} \\
 & + 3N\delta_{a_1b_4}\delta_{a_2b_1}\delta_{a_3b_3} + N\delta_{a_1b_3}\delta_{a_3b_1}\delta_{a_2b_4} - 4\delta_{a_1b_1}\delta_{a_3b_3}\delta_{a_2b_4}] \tag{15}
 \end{aligned}$$

Making use of Eq.(3) one has

$$\begin{aligned}
 & \beta_{kln}\beta_{nmo}\beta_{lmp}(\lambda_k)_{a_1b_1}(\lambda_p)_{a_2b_2}(\lambda_o)_{a_4b_4} \\
 = & \alpha^3[-2N^2\delta_{a_1b_3}\delta_{a_3b_4}\delta_{a_2b_1} - N^2\delta_{a_1b_4}\delta_{a_2b_3}\delta_{a_3b_1} + 3N^2\delta_{a_1b_1}\delta_{a_2b_3}\delta_{a_3b_4} \\
 & + 3N\delta_{a_1b_4}\delta_{a_2b_1}\delta_{a_3b_3} + N\delta_{a_1b_3}\delta_{a_3b_1}\delta_{a_2b_4} - 4\delta_{a_1b_1}\delta_{a_3b_3}\delta_{a_2b_4}] \\
 & - \alpha^2(N^2 - 2)(\lambda_k)_{a_1b_1}(\lambda_k)_{a_4b_4}\delta_{a_2b_3} \tag{16}
 \end{aligned}$$

here we note that N is arbitrary integer, let $N = 3$, $\alpha = \frac{2}{3}$. That is, for the $SU(3)$ case, with Eq.(10), one has

$$\delta_{b_1b_2}\delta_{a_2a_3} = \frac{1}{2}(\lambda_k)_{b_1a_3}(\lambda_k)_{a_2b_2} + \frac{1}{3}\delta_{b_1a_3}\delta_{a_2b_2} \tag{17}$$

Consider $\delta_{b_1b_2}\delta_{a_2a_3}\delta_{b_3a_1}$:

$$\begin{aligned}
 & \delta_{b_1b_2}\delta_{a_2a_3}\delta_{b_3a_1} \\
 = & \left[\frac{1}{2}(\lambda_k)_{b_1a_3}(\lambda_k)_{a_2b_2} + \frac{1}{3}\delta_{b_1a_3}\delta_{a_2b_2}\right]\delta_{b_3a_1} \\
 = & \frac{1}{4}(\lambda_k\lambda_l)_{b_1a_1}(\lambda_k)_{a_2b_2}(\lambda_l)_{b_3a_3} + \frac{1}{6}(\lambda_k)_{b_1a_1}(\lambda_k)_{a_2b_2}\delta_{b_3a_3} \\
 & + \frac{1}{6}(\lambda_k)_{b_1a_1}(\lambda_k)_{b_3a_3}\delta_{a_2b_2} + \frac{1}{9}\delta_{b_1a_1}\delta_{a_2b_2}\delta_{b_3a_3} \tag{18}
 \end{aligned}$$

where one has made use of

$$\begin{aligned}
 (\lambda_k)_{b_1a_3}\delta_{b_3a_1} &= (\lambda_k)_{b_1a_4}\delta_{a_4a_3}\delta_{b_3a_1} \\
 &= (\lambda_k)_{b_1a_4}\left[\frac{1}{2}(\lambda_l)_{a_4a_1}(\lambda_l)_{b_3a_3} + \frac{1}{3}\delta_{a_4a_1}\delta_{b_3a_3}\right] \\
 &= \frac{1}{2}(\lambda_k\lambda_l)_{b_1a_1}(\lambda_l)_{b_3a_3} + \frac{1}{3}(\lambda_k)_{b_1a_1}\delta_{b_3a_3} \tag{19}
 \end{aligned}$$

Substituting Eqs.(17,18) into Eq.(16), one has

$$\begin{aligned}
& \beta_{kln}\beta_{nmo}\beta_{lmp}(\lambda_o)_{b_3a_3}(\lambda_p)_{b_1a_1}(\lambda_k)_{a_2b_2} \\
= & \frac{8}{27}[30\delta_{a_2a_3}\delta_{b_1a_1}\delta_{b_3b_2} - 18\delta_{a_2a_1}\delta_{b_1a_3}\delta_{b_3b_2} - 13\delta_{a_2b_2}\delta_{b_1a_1}\delta_{b_3a_3} \\
& - 9\delta_{a_2a_3}\delta_{b_1b_2}\delta_{b_3a_1} + 9\delta_{a_2b_2}\delta_{b_1a_3}\delta_{b_3a_2} + 9\delta_{a_2a_1}\delta_{b_1b_2}\delta_{b_3a_3}] - \frac{28}{9}(\lambda_k)_{a_2b_2}(\lambda_k)_{b_3a_3}\delta_{b_1a_1} \\
= & \frac{8}{27}[30(\frac{1}{2}(\lambda_k)_{a_2b_2}(\lambda_k)_{b_3a_3} + \frac{1}{3}\delta_{a_2b_2}\delta_{b_3a_3})\delta_{b_1a_1} - 18[\frac{1}{4}(\lambda_k)_{b_1a_1}(\lambda_k\lambda_m)_{a_2b_2}(\lambda_m)_{b_3a_3} \\
& + \frac{1}{6}(\lambda_k)_{b_1a_1}(\lambda_k)_{a_2b_2}\delta_{a_3b_3} + \frac{1}{6}(\lambda_k)_{a_2b_2}(\lambda_k)_{b_3a_3}\delta_{b_1a_1} + \frac{1}{9}\delta_{b_1a_1}\delta_{a_2b_2}\delta_{b_3a_3}] \\
& - 13\delta_{b_3a_3}\delta_{b_1a_1}\delta_{a_2b_2} - 9[\frac{1}{4}(\lambda_k\lambda_m)_{b_1a_1}(\lambda_k)_{a_2b_2}(\lambda_m)_{b_3a_3} + \frac{1}{6}(\lambda_k)_{b_1a_1}(\lambda_k)_{a_2b_2}\delta_{b_3a_3} \\
& + \frac{1}{6}(\lambda_k)_{b_1a_1}(\lambda_k)_{a_2b_2}\delta_{b_3a_3} + \frac{1}{6}(\lambda_k)_{b_1a_1}(\lambda_k)_{b_3a_3}\delta_{a_2b_2} + \frac{1}{9}\delta_{b_1a_1}\delta_{a_2b_2}\delta_{a_3b_3}] \\
& + 9[\frac{1}{2}(\lambda_k)_{b_1a_1}(\lambda_k)_{b_3a_3} + \frac{1}{3}\delta_{b_1a_1}\delta_{b_3a_3}]\delta_{a_2b_2} + 9[\frac{1}{2}(\lambda_k)_{a_1b_2}(\lambda_k)_{b_1a_1} \\
& + \frac{1}{3}\delta_{b_1a_1}\delta_{a_2b_2}]\delta_{b_3a_3}) - \frac{28}{9}(\lambda_k)_{a_2b_2}(\lambda_k)_{b_3a_3}\delta_{b_1a_1} \\
= & -\frac{4}{3}\beta_{klm}^*(\lambda_k)_{b_3a_3}(\lambda_l)_{b_1a_1}(\lambda_m)_{a_2b_2} - \frac{2}{3}\beta_{klm}(\lambda_k)_{b_3a_3}(\lambda_l)_{b_1a_1}(\lambda_m)_{a_2b_2} \tag{20}
\end{aligned}$$

Similarly, one can obtain

$$\begin{aligned}
\beta_{klp}\beta_{kmn}\beta_{lmo}(\lambda_p)_{b_3a_3}(\lambda_n)_{b_1a_1}(\lambda_m)_{a_2b_2} = & -\frac{4}{3}\beta_{k'm}(\lambda_k)_{b_3a_3}(\lambda_l)_{b_1a_1}(\lambda_m)_{a_2b_2} \\
& -\frac{2}{3}\beta_{klm}^*(\lambda_k)_{b_3a_3}(\lambda_l)_{b_1a_1}(\lambda_m)_{a_2b_2} \tag{21}
\end{aligned}$$

$$\beta_{kln}\beta_{nmo}\beta_{mlp}(\lambda_o)_{b_3a_3}(\lambda_k)_{b_1a_1}(\lambda_p)_{a_2b_2} = 4\beta_{klm}(\lambda_k)_{b_3a_3}(\lambda_l)_{b_1a_1}(\lambda_m)_{a_2b_2} \tag{22}$$

These relations have been checked numerically by a computer program.

3 DISCUSSION

These relations (Eqs.(20, 21 and 22)) may be useful in deriving many body interaction. For example, when we derive the $qq\bar{q}-qq\bar{q}$ interaction in the QCD^[2], one needs Eq.(21).

The colour matrix element of the $q\bar{q}q$ - $q\bar{q}q$ interaction can be written as:

$$\begin{aligned}
 & (\lambda^k \lambda^l)_{c'_1 c_1} (\lambda^k)_{c'_2 c'_3} (\lambda^l)_{c_3 c_2} \\
 &= \chi_{c'_1} \chi_{c'_2} \chi_{c_3} \left\{ \left(\frac{2}{3} \delta_{kl} + \beta_{klp} \lambda_1^p \right) \left[\frac{1}{3} \lambda_2^k \lambda_3^l + \beta_{kmn} \beta_{kmo} \lambda_2^n \lambda_3^o \right. \right. \\
 & \quad \left. \left. + \frac{2}{9} \delta_{kl} + \frac{1}{3} \beta_{kln} \lambda_2^n + \frac{1}{3} \beta_{lko} \lambda_3^o \right] \right\} \chi_{c_1} \chi_{c_2} \chi_{c'_3} \\
 &= \chi_{c'_1} \chi_{c'_2} \chi_{c_3} \left\{ -\frac{2}{9} \vec{\lambda}_2 \cdot \vec{\lambda}_3 - \frac{4}{9} \vec{\lambda}_1 \cdot \vec{\lambda}_2 + \frac{4}{9} \vec{\lambda}_1 \cdot \vec{\lambda}_3 + \frac{32}{27} \right. \\
 & \quad \left. + \frac{1}{3} \beta_{klp} \lambda_1^p \lambda_2^k \lambda_3^l + \frac{1}{2} \beta_{klp} \beta_{kmn} \beta_{lmo} \lambda_1^p \lambda_2^n \lambda_3^o \right\} \\
 &= \chi_{c'_1} \chi_{c'_2} \chi_{c_3} \left(\frac{1}{9} (\lambda_2 - \lambda_3^*)^2 + \frac{16}{9} \lambda_1 \cdot \lambda_3^* - \frac{2}{9} \lambda_1 \cdot \lambda_2 \right. \\
 & \quad \left. - \frac{1}{3} \lambda_1 \cdot \lambda_3^* \lambda_2 \cdot \lambda_3^* - \frac{1}{3} \lambda_2 \cdot \lambda_3^* \lambda_1 \cdot \lambda_2 \right) \chi_{c_1} \chi_{c_2} \chi_{c_3} \tag{23}
 \end{aligned}$$

which has been made use of the generalized Fierz transformation and Eq.(21).

REFERENCES

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